

PHYS 452: Quantum Mechanics II, Quiz #4**Instruction: use additional sheets if you find it necessary**

A particle is incident on a central potential $V(r)$, which is infinitely high at $r \leq a$ and vanish at $r > a$, i.e.

$$V(r) = \begin{cases} 0, & r > a \\ \infty, & r \leq a \end{cases}$$

Find the total cross section when the energy of the incident particle is low. Define what “low” means in this context.

Appendix:**Stationary quantum scattering**

Wave function at $r \rightarrow \infty$: $\psi(r, \theta, \phi) \approx A \left[e^{ikz} + f(\theta, \phi) \frac{e^{ikr}}{r} \right]$, $k = \frac{\sqrt{2mE}}{\hbar}$

Differential cross section: $\frac{d\sigma}{d\Omega} = |f(\theta, \phi)|^2$ Total cross section: $\sigma_{\text{tot}} = \int \frac{d\sigma}{d\Omega} d\Omega$

Partial wave analysis

For a spherically symmetric potentials $\psi(r, \theta) = A \left[e^{ikz} + k \sum_{l=0}^{\infty} i^{l+1} (2l+1) a_l h_l^{(1)}(kr) P_l(\cos \theta) \right]$

$$f(\theta) = \sum_{l=0}^{\infty} (2l+1) a_l P_l(\cos \theta) = \frac{1}{k} \sum_{l=0}^{\infty} (2l+1) e^{i\delta_l} \sin \delta_l P_l(\cos \theta)$$

$$\sigma_{\text{tot}} = 4\pi \sum_{l=0}^{\infty} (2l+1) |a_l|^2 = \frac{4\pi}{k^2} \sum_{l=0}^{\infty} (2l+1) \sin^2 \delta_l$$

Relation between partial wave amplitudes and phase shifts: $a_l = \frac{1}{k} e^{i\delta_l} \sin \delta_l$

Rayleigh formula for a plane wave expansion: $e^{ikz} = \sum_{l=0}^{\infty} i^l (2l+1) j_l(kr) P_l(\cos \theta)$

Lippmann-Schwinger equation

$$\psi(\mathbf{r}) = \varphi(\mathbf{r}) + \frac{2m}{\hbar^2} \int G(\mathbf{r}, \mathbf{r}') V(\mathbf{r}') \psi(\mathbf{r}') d\mathbf{r}'$$

where $\varphi(\mathbf{r})$ is the free-particle solution (incident plane wave)

and $G(\mathbf{r}, \mathbf{r}') = -\frac{1}{4\pi} \frac{e^{ik|\mathbf{r}-\mathbf{r}'|}}{|\mathbf{r}-\mathbf{r}'|}$ is the Green's function

Born approximation

$$f(\theta, \phi) = -\frac{m}{2\pi\hbar^2} \int e^{i\mathbf{q}\cdot\mathbf{r}'} V(\mathbf{r}') d\mathbf{r}', \quad \mathbf{q} = \mathbf{k}' - \mathbf{k}, \quad q = 2k \sin \frac{\theta}{2}, \quad \mathbf{k} = k\hat{\mathbf{r}}, \quad \mathbf{k}' = k\hat{\mathbf{z}}$$

For spherically symmetric potentials $f(\theta) = -\frac{2m}{\hbar^2 q} \int_0^{\infty} r V(r) \sin(qr) dr$

Legendre polynomials

$$P_0(x) = 1, \quad P_1(x) = x, \quad P_2(x) = \frac{3}{2}x^2 - \frac{1}{2}, \quad P_3(x) = \frac{5}{2}x^3 - \frac{3}{2}x, \quad \dots, \quad P_l(x) = \frac{1}{2^l l!} \left(\frac{d}{dx} \right)^l (x^2 - 1)^l$$

$$\text{Orthogonality: } \int_{-1}^1 P_l(x) P_{l'}(x) dx = \frac{2}{2l+1} \delta_{ll'}$$

Spherical Bessel, Neumann, and Hankel functions

$$\begin{aligned}
j_0(x) &= \frac{\sin x}{x}, & j_1(x) &= \frac{\sin x}{x^2} - \frac{\cos x}{x}, & \dots, & & j_l(x) &= (-x)^l \left(\frac{1}{x} \frac{d}{dx}\right)^l \frac{\sin x}{x} \\
n_0(x) &= -\frac{\cos x}{x}, & n_1(x) &= -\frac{\cos x}{x^2} - \frac{\sin x}{x}, & \dots, & & n_l(x) &= -(-x)^l \left(\frac{1}{x} \frac{d}{dx}\right)^l \frac{\cos x}{x} \\
h_l^{(1,2)}(x) &= j_l(x) \pm i n_l(x) \\
h_0^{(1)}(x) &= -i \frac{e^{ix}}{x}, & h_1^{(1)}(x) &= \left(-\frac{i}{x^2} - \frac{1}{x}\right) e^{ix}, & h_2^{(1)}(x) &= \left(-\frac{3i}{x^3} - \frac{3}{x^2} + \frac{i}{x}\right) e^{ix}, & \dots \\
h_0^{(2)}(x) &= i \frac{e^{-ix}}{x}, & h_1^{(2)}(x) &= \left(\frac{i}{x^2} - \frac{1}{x}\right) e^{-ix}, & h_2^{(2)}(x) &= \left(\frac{3i}{x^3} - \frac{3}{x^2} + \frac{i}{x}\right) e^{-ix}, & \dots \\
\text{For } x \ll 1: & j_l(x) \rightarrow \frac{2^l l!}{(2l+1)!} x^l, & n_l &\rightarrow -\frac{(2l)!}{2^l l!} x^{-l-1} \\
\text{For } x \gg 1: & h_l^{(1)} \rightarrow \frac{1}{x} (-i)^{l+1} e^{ix}, & h_l^{(2)} &\rightarrow \frac{1}{x} (i)^{l+1} e^{-ix}
\end{aligned}$$

Useful integrals

$$\begin{aligned}
\int_0^{\infty} x^{2k} e^{-\beta x^2} dx &= \sqrt{\pi} \frac{(2k)!}{k! 2^{2k+1} \beta^{k+1/2}} \quad (\text{Re } \beta > 0, k = 0, 1, 2, \dots) \\
\int_0^{\infty} x^{2k+1} e^{-\beta x^2} dx &= \frac{1}{2} \frac{k!}{\beta^{k+1}} \quad (\text{Re } \beta > 0, k = 0, 1, 2, \dots) \\
\int_0^{\infty} x^k e^{-\gamma x} dx &= \frac{k!}{\gamma^{k+1}} \quad (\text{Re } \gamma > 0, k = 0, 1, 2, \dots) \\
\int_{-\infty}^{\infty} e^{-\beta x^2} e^{iqx} dx &= \sqrt{\frac{\pi}{\beta}} e^{-\frac{q^2}{4\beta}} \quad (\text{Re } \beta > 0) \\
\int_0^{\pi} \sin^{2k} x dx &= \pi \frac{(2k-1)!!}{2^k k!} \quad (k = 0, 1, 2, \dots) \\
\int_0^{\pi} \sin^{2k+1} x dx &= \frac{2^{k+1} k!}{(2k+1)!!} \quad (k = 0, 1, 2, \dots)
\end{aligned}$$

Useful trigonometric identities

$$\begin{aligned}
\sin(\alpha \pm \beta) &= \sin \alpha \cos \beta \pm \cos \alpha \sin \beta & \cos(\alpha \pm \beta) &= \cos \alpha \cos \beta \mp \sin \alpha \sin \beta \\
\sin \alpha \sin \beta &= \frac{1}{2} [\cos(\alpha - \beta) - \cos(\alpha + \beta)] & \cos \alpha \cos \beta &= \frac{1}{2} [\cos(\alpha - \beta) + \cos(\alpha + \beta)] \\
\sin \alpha \cos \beta &= \frac{1}{2} [\sin(\alpha + \beta) + \sin(\alpha - \beta)] & \cos \alpha \sin \beta &= \frac{1}{2} [\sin(\alpha + \beta) - \sin(\alpha - \beta)]
\end{aligned}$$